

# Interface Debonding Prediction Technique for Tensile Loaded AA6061/Zirconium Oxide Nanoparticulate MMC

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**Abstract:** A diamond array unit cell/2-D octahedron particulate RVE models were employed to evaluate interfacial debonding using cohesive zone analysis. The particulate metal matrix composites are zirconium oxide/AA6061 alloy at volume fractions of 10%, 20% and 30% boron nitride. Interface debonding was observed in all the composites at angle of 72° from axis of loading.

**Keywords:** AA6061 alloy, zirconium oxide, octahedron particulate, RVE model, finite element analysis, interface debonding.

## 1. INTRODUCTION

Metal matrix composite (MMC) materials are superior to conventional monolithic materials for a number of reasons, noteworthy amongst these being their mechanical properties. The mechanism of failure at the microscale is dominated by factors such as the cohesive strength of the interface, and the strength and reliability of the particle and matrix phase. The interface between the reinforcement and matrix plays a crucial role in changing the property of composite material. In many cases the properties of MMC can be drastically altered by the nature of the interface, making it necessary to understand the characteristics of the interface. A number of different approaches have been considered for modeling the interface, among which the cohesive zone models are widely used to model delamination and debonding in composite materials [1-14].

The present study is on the interface debonding between AA6061 alloy and zirconium oxide nanoparticle. Representative volume elements (RVEs) models were taken from the periodic 2-D octahedron particulates in a diamond array distribution (figure 1).

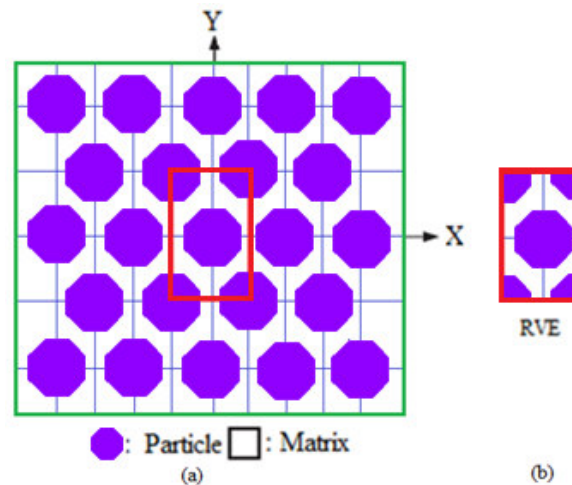


Figure 1: The RVE model.

## 2. MATERIALS AND METHODS

The computational domain considered in the current research is comprised of AA6061 alloy matrix material with an embedded 2-D octahedron zirconium oxide inclusion. The volume fractions of zirconium oxide were 10%, 20%, and 30%. In the model periodic boundary conditions were used in the x and y directions. Initially, both AA6061 and zirconium oxide were kept in contact along the interface in the x-y plane with zero separation distance. The simulation domain is divided into three regions: zirconium oxide nanoparticle, interface and AA6061 alloy matrix. PLANE183 element was used for the matrix and the

nanoparticulates. The cohesive zone can be incorporated in the continuum formulation by applying the cohesive tractions as boundary conditions. The cohesive element is implemented as a linear element with four nodes. Initially, the interface between the matrix material and the inclusion is assumed to be perfectly bonded, that is, continuity of traction and displacement is assumed along the interface. The finite element analysis was carried out for the single inclusion model undergoing a tensile load. The elastic material properties are given by  $E_m = 68.90$  GPa,  $E_p = 250$  GPa,  $\nu_m = 0.33$  and  $\nu_p = 0.32$ .

Shear-lag model is based on the assumption that all of the load transfer from matrix to particulate occurs via shear stresses acting on the particulate interface between the two constituents. The rate of change of the stress in the particulate to the interfacial shear stress at that point and the particulate radius, 'r' is given by:

$$\frac{d\sigma_p}{dx} = -\frac{2\tau_i}{r} \quad (1)$$

which may be regarded as the basic shear lag relationship.

The stress distribution in the particulate is determined by relating shear strains in the matrix around the particulate to the macroscopic strain of the composite. Some mathematical manipulation leads to a solution for the distribution of stress at a distance 'x' from the mid-point of the particulate which involves hyperbolic trig functions:

$$\sigma_p = E_p \varepsilon_c [1 - \cosh(nx/r) \operatorname{sech}(ns)] \quad (2)$$

where  $\varepsilon_c$  is the composite strain,  $s$  is the particulate aspect ratio (length/diameter) and  $n$  is a dimensionless constant given by:

$$n = \left[ \frac{2E_m}{E_p(1+\nu_m)\ln(1+\nu_p)} \right]^{1/2} \quad (3)$$

in which  $\nu_m$  is the Poisson ratio of the matrix. The variation of interfacial shear stress along the particulate length is derived, according to Equation (1), by differentiating this equation, to give:

$$\tau_i = \frac{n\varepsilon_c}{2} E_p \sinh\left(\frac{nx}{r}\right) \operatorname{sech}(ns) \quad (4)$$

The equation for the stress in the particulate, together with the assumption of a average tensile strain in the matrix equal to that imposed on the composite, can be used to evaluate the composite stiffness. This leads to:

$$\sigma_c = \varepsilon_c \left[ \nu_p E_p \left(1 - \frac{\tanh(ns)}{ns}\right) + (1 - \nu_p) E_m \right] \quad (5)$$

The expression in square brackets is the composite stiffness. The stiffness is a function of particulate aspect ratio, particulate/matrix stiffness ratio and particulate volume fraction.

If the particle deforms in an elastic manner (according to Hooke's law) then,

$$\tau = \frac{n}{2} \sigma_p \quad (6)$$

If interfacial debonding/yielding is considered to occur when the interfacial shear stress reaches its shear strength

$$\tau = \tau_{\max} \quad (7)$$

For particle/matrix interfacial fracture can be established whereby,

$$\tau_{\max} < \frac{n\sigma_p}{2} \quad (8)$$

This approach suggests that the outcome of a matrix crack impinging on an embedded particle depends on the balance between the particle strength and the shear strength of the interface. For plane strain conditions, the macro stress- macro strain relation is as follows:

$$\begin{Bmatrix} \overline{\sigma_x} \\ \overline{\sigma_y} \\ \overline{\tau_{xy}} \end{Bmatrix} = \begin{bmatrix} \overline{C_{11}} & \overline{C_{12}} & 0 \\ \overline{C_{21}} & \overline{C_{22}} & 0 \\ 0 & 0 & \overline{C_{33}} \end{bmatrix} \times \begin{Bmatrix} \overline{\varepsilon_x} \\ \overline{\varepsilon_y} \\ \overline{\gamma_{xy}} \end{Bmatrix} \quad (9)$$

The interfacial tractions can be obtained by transforming the micro stresses at the interface as given in Eq. (3):

$$t = \begin{Bmatrix} t_z \\ t_n \\ t_t \end{Bmatrix} = T\sigma \quad (10)$$

$$\text{where, } T = \begin{bmatrix} 0 & 0 & 0 \\ \cos^2\theta & \sin^2\theta & 2\sin\theta\cos\theta \\ -\sin\theta\cos\theta & \sin\theta\cos\theta & \cos^2\theta - \sin^2\theta \end{bmatrix}$$

### 3. RESULTS AND DISCUSSION

The tensile and compression moduli increased with volume fraction of  $ZrO_2$  as shown figure 2a. The shear modulus increased constant with increase in the volume fraction of  $ZrO_2$  up to 20%  $V_p$  in the composites (figure 2b), while it was lower for 30% $V_p$  of  $ZrO_2$  than that of 10% and 20% $V_p$  of  $ZrO_2$ . The major Poisson's ratio decreased with volume fraction of  $ZrO_2$ . The

stiffness mismatch between ZrO<sub>2</sub> nanoparticulate and AA6061 alloy matrix is 181 GPa. The condition  $\tau_{max} < n\sigma_p/2$  is satisfied for the prevalence of interface debonding in the composites including 10%, 20% and 30% ZrO<sub>2</sub> (figure 3). The strain energy density developed in the composites is shown in figures 4 and 5. For the shearing of interface between ZrO<sub>2</sub> inclusion and AA6061 alloy matrix, the strain energy density at the interface was higher than that in ZrO<sub>2</sub> particle and AA6061 alloy matrix.

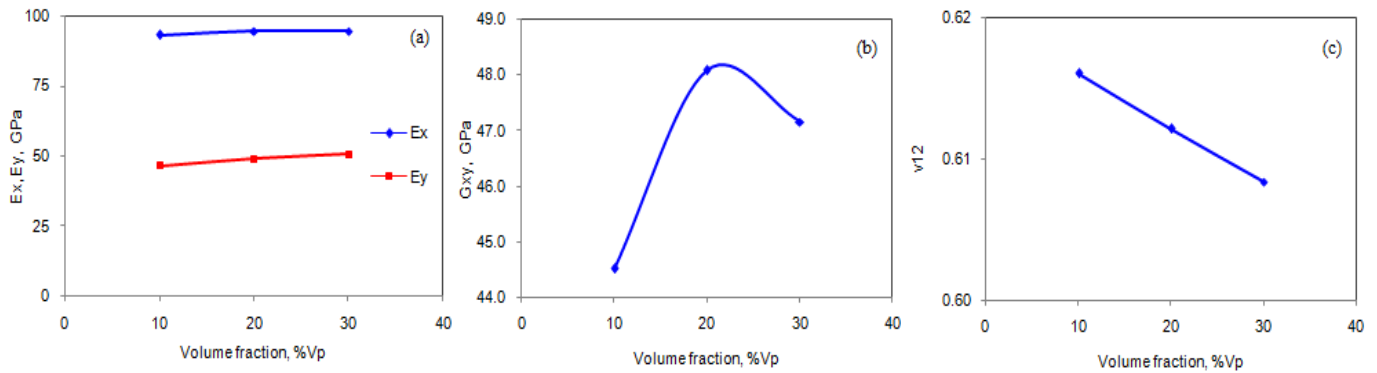


Figure 2: Effect of volume fraction on effective material properties.

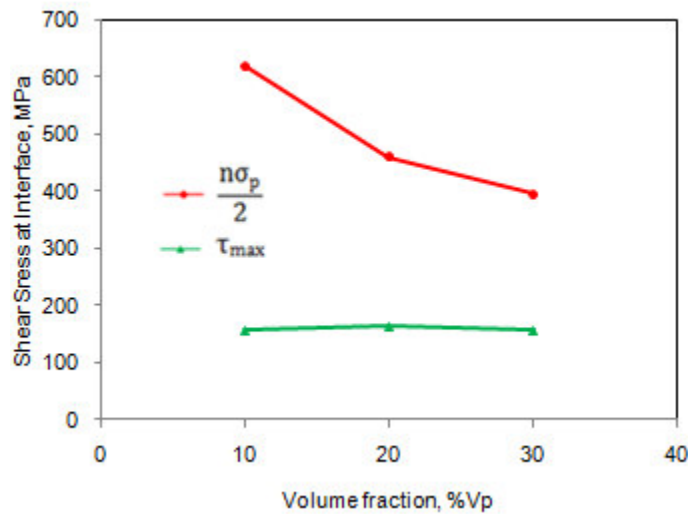


Figure 3: Fracture criteria of interface debonding.

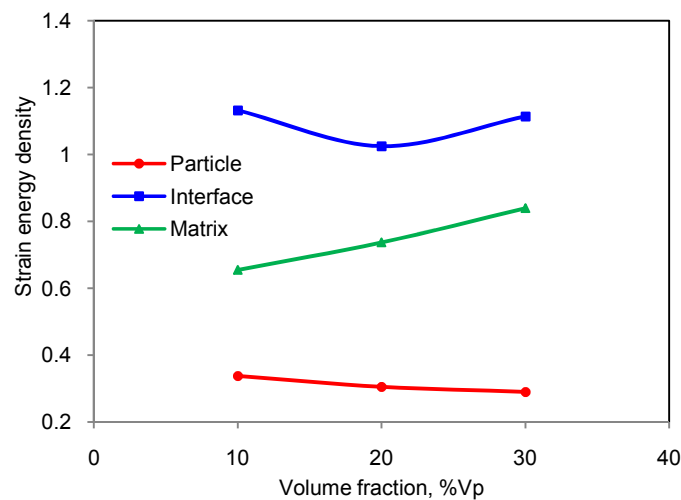


Figure 4: Effect of volume fraction on strain energy density.

The normal and tangential tractions are plotted in figure 6. As these traction stresses were calculated locally near the interface. The normal traction in the region of interface between  $ZrO_2$  nanoparticulate and AA6061 alloy matrix was high in the direction ( $0^\circ$ ) of tensile loading while it decreased with increase of angle along the interface. The maximum allowable normal and tangential separation of the cohesive element corresponds to the separation distance at which the traction decreases to near zero value. The separation took place at  $72^\circ$  from the axis of loading for the octahedron nanoparticulate in the matrix of AA6061 alloy. The strain energy density at point of separation was 0.427 for the composites having 30%  $ZrO_2$ .

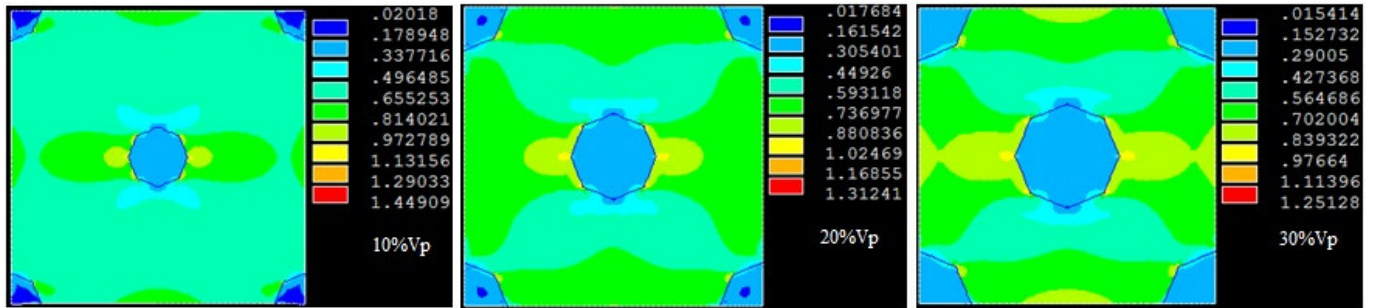


Figure 5: FEA results of strain energy densities.

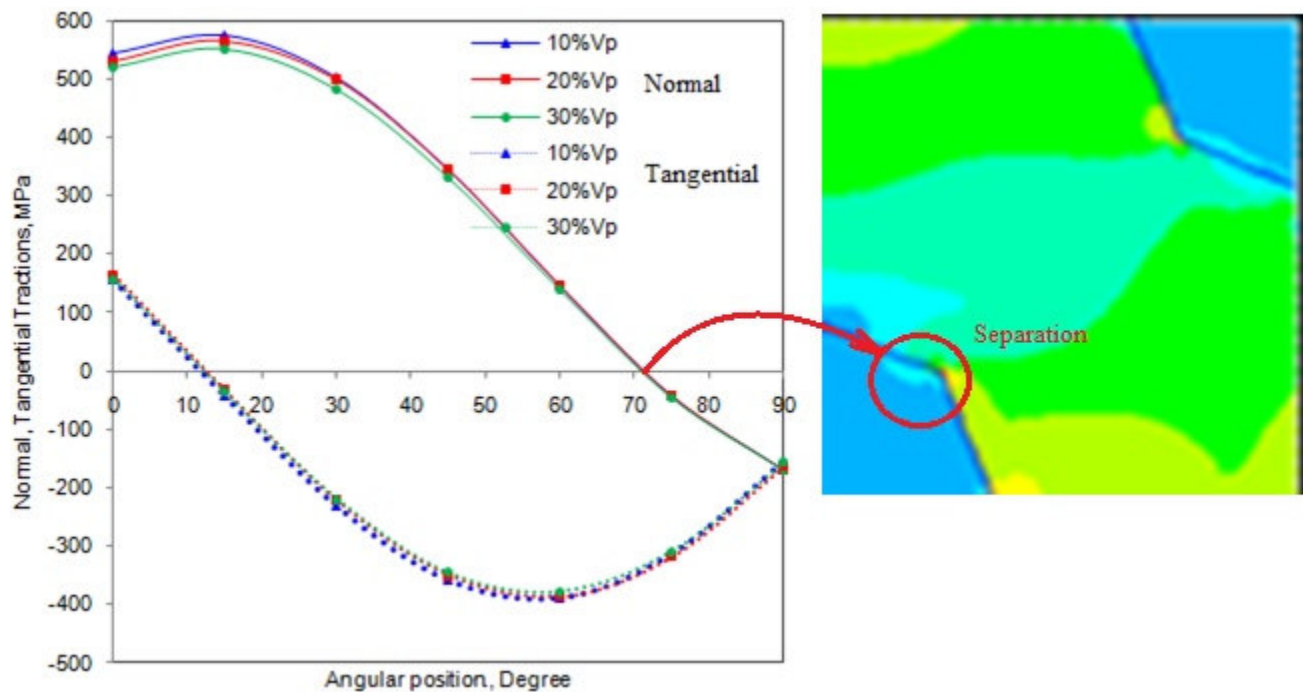


Figure 6: Normal and tangential: (a) tractions and (b) displacements.

#### 4. CONCLUSION

The interface debonding occurred in the composites containing 10%, 20% and 30% volume fractions  $ZrO_2$ . As these stresses were calculated locally near the interface they therefore relate to the traction in the traction–separation law. The maximum strain densities at which the traction decreases to near zero value were 0.497, 0.449 and 0.427, respectively, for 10%, 20% and 30%  $ZrO_2$  inclusions in AA6061 alloy matrix.

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